

Representations of Lie Algebras and Systems on Lie Groups

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Lie Algebra Operator Vessels

- ▶ We fix a finite dimensional real Lie algebra \mathfrak{g} and a linear map $\rho: \mathfrak{g} \rightarrow \mathcal{L}(\mathcal{H})$, for a Hilbert space \mathcal{H} , such that $\frac{1}{i}\rho$ is a representation.
- ▶ A \mathfrak{g} -operator vessel is a tuple:

$$\mathfrak{V} = (\mathcal{H}, \mathcal{E}, \rho, \Phi, \sigma, \gamma, \gamma_*).$$

- ▶ Here \mathcal{E} is the auxiliary Hilbert space, $\Phi \in \mathcal{L}(\mathcal{H}, \mathcal{E})$, σ is a linear map $\sigma: \mathfrak{g} \rightarrow \mathcal{L}(\mathcal{E})$ and γ and γ_* are linear maps $\gamma, \gamma_*: \mathfrak{g} \wedge \mathfrak{g} \rightarrow \mathcal{L}(\mathcal{E})$.

Lie Algebra Operator Vessels

- ▶ Furthermore the tuple satisfies the following conditions:
 - ▶ $\forall X, Y \in \mathfrak{g}: \gamma(X \wedge Y) - \gamma(X \wedge Y)^* = i\sigma([X, Y]),$
 - ▶ $\forall X, Y \in \mathfrak{g}: \gamma_*(X \wedge Y) - \gamma_*(X \wedge Y)^* = i\sigma([X, Y]),$
 - ▶ The colligation condition: $\forall X \in \mathfrak{g}: \frac{1}{i}(\rho(X) - \rho(X)^*) = \Phi^* \sigma \Phi,$
 - ▶ The input condition:
 $\forall X, Y \in \mathfrak{g}: \sigma(X)\Phi\rho(Y)^* - \sigma(Y)\Phi\rho(X)^* = \gamma(X \wedge Y)\Phi.$
 - ▶ The output condition:
 $\forall X, Y \in \mathfrak{g}: \sigma(X)\Phi\rho(Y) - \sigma(Y)\Phi\rho(X) = \gamma_*(X \wedge Y)\Phi,$
 - ▶ The linkage condition: $\forall X, Y \in \mathfrak{g}: \gamma_*(X \wedge Y) - \gamma(X \wedge Y) = i(\sigma(X)\Phi\Phi^*\sigma(Y) - \sigma(Y)\Phi\Phi^*\sigma(X)).$

Lie Algebra Operator Vessels

- ▶ Let us consider a simple example. If $\mathfrak{g} = \mathbb{R}^n$. Denote by X_1, \dots, X_n the canonical basis. Set $A_j = \rho(X_j)$ for every $j = 1, \dots, n$.
- ▶ Since $[X, Y] = 0$ for every $X, Y \in \mathfrak{g}$, from the first two vessel conditions we get that $\gamma(X) = \gamma(X)^*$ and $\gamma_*(X) = \gamma_*(X)^*$, for every $X \in \mathfrak{g}$.
- ▶ We thus recover a commutative n -operator vessel as defined in [Livsic et al, 1995]

Lie Algebra Operator Vessels

- ▶ The $ax + b$ is the only (up to isomorphism) non-commutative Lie group of the dimension 2.
- ▶ The $ax + b$ group can be realised as a subgroup of $GL_2(\mathbb{R})$, of matrices of the form:

$$\begin{pmatrix} a & b \\ 0 & 1 \end{pmatrix}, a, b \in \mathbb{R}, a > 0$$

- ▶ The $ax + b$ associated Lie algebra can be realized as a subalgebra of $M_2(\mathbb{R})$ of the form:

$$\begin{pmatrix} p & q \\ 0 & 0 \end{pmatrix}, p, q \in \mathbb{R},$$

Lie Algebra Operator Vessels

- ▶ The fields $X_1 = a \frac{\partial}{\partial a}$ and $X_2 = a \frac{\partial}{\partial b}$ span the algebra.
- ▶ Let us make the identification $A_j = \rho(X_j)$, $\sigma_j = \sigma(X_j)$, $\gamma = \gamma(X_1 \wedge X_2)$ and $\gamma_* = \gamma_*(X_1 \wedge X_2)$.
- ▶ Then we have $[A_1, A_2] = \frac{1}{i} A_2$ and $\gamma - \gamma_* = i\sigma_2$.
- ▶ We will carry this example on and use the same notations.

Lie Algebra Operator Vessels

- ▶ Every Lie algebra representation, $\tau: \mathfrak{g} \rightarrow \mathcal{L}(\mathcal{H})$, can be embedded in a Lie algebra operator vessel.
- ▶ Set $\rho = i\tau$, $\mathcal{E} = \overline{\sum_{X \in \mathfrak{g}} (\rho(X) - \rho(X)^*)\mathcal{H}}$ and $\Phi = P_{\mathcal{E}}$, the orthogonal projection onto \mathcal{E} .
- ▶ We then define:
 - ▶ $\sigma(X) = \frac{1}{i}(\rho(X) - \rho(X)^*)|_{\mathcal{E}}$,
 - ▶ $\gamma(X \wedge Y) = \frac{1}{i}((\rho(X)\rho(Y)^* - \rho(Y)\rho(X)^*) - \rho([X, Y]^*))|_{\mathcal{E}}$,
 - ▶ $\gamma_*(X \wedge Y) = \frac{1}{i}((\rho(Y)^*\rho(X) - \rho(X)^*\rho(Y)) + \rho([X, Y]))|_{\mathcal{E}}$.
- ▶ Note that in the case $\mathfrak{g} = \mathbb{R}^n$ we recover the construction of a vessel for an n -tuple of commuting operators.

Lie Algebra Operator Vessels

- ▶ Let \mathcal{G} be a connected, simply connected Lie group associated to \mathfrak{g} .
- ▶ The vessel defines a linear, time invariant, overdetermined system on \mathcal{G} :

$$\begin{cases} iXf + \rho(X)f = \Phi^* \sigma(X)u, \\ y = u - i\Phi f \end{cases}$$

- ▶ Here $X \in \mathfrak{g}$ and $u, y: \mathcal{G} \rightarrow \mathcal{E}$ and $f: \mathcal{G} \rightarrow \mathcal{H}$ are smooth functions. We call u the input of the system, y the output and f the state.

Lie Algebra Operator Vessels

- ▶ Since the system is overdetermined, it admits a set of compatibility conditions in the input:

$$\forall X, Y \in \mathfrak{g}: \sigma(Y)Xu - \sigma(X)Yu + i\gamma(X \wedge Y)u = 0.$$

- ▶ One can show that if $u: \mathfrak{G} \rightarrow \mathcal{E}$ is a function satisfying the compatibility equations, then there exists $f: \mathfrak{G} \rightarrow \mathcal{H}$ that solves the system.
- ▶ Furthermore the output thus defined satisfies the output compatibility equations:

$$\forall X, Y \in \mathfrak{g}: \sigma(Y)Xy - \sigma(X)Yy + i\gamma_*(X \wedge Y)y = 0.$$

Lie Algebra Operator Vessels

- ▶ Let $\mathfrak{g} = \mathbb{R}^2$. Then $\mathfrak{G} = \mathbb{R}^2$ as well.
- ▶ Let $\sigma_1 = \sigma(X_1)$, $\sigma_2 = \sigma(X_2)$, $\gamma = \gamma(X_1 \wedge X_2)$ and $\gamma_* = \gamma_*(X_1 \wedge X_2)$.
- ▶ The system equations are:

$$\begin{cases} i \frac{\partial f}{\partial t_j} + A_j f = \Phi^* \sigma_j u, \\ y = u - i\Phi f \end{cases}$$

Lie Algebra Operator Vessels

- ▶ The input compatibility equation is then:

$$\sigma_2 \frac{\partial u}{\partial t_1} - \sigma_1 \frac{\partial u}{\partial t_2} + i\gamma u = 0.$$

- ▶ The output compatibility equation is:

$$\sigma_2 \frac{\partial u}{\partial t_1} - \sigma_1 \frac{\partial u}{\partial t_2} + i\gamma_* u = 0.$$

Lie Algebra Operator Vessels

- ▶ The $ax + b$ system of equations has the following form:

$$ia \frac{\partial f}{\partial a} + A_1 f = \Phi^* \sigma_1 u,$$

$$ia \frac{\partial f}{\partial b} + A_2 f = \Phi^* \sigma_2 u,$$

$$y = u - i\Phi f.$$

- ▶ The input compatibility equations are:

$$a\sigma_2 \frac{\partial u}{\partial a} - a\sigma_1 \frac{\partial u}{\partial b} + i\gamma u = 0.$$

Lie Group Representations

- ▶ Let \mathfrak{G} be a Lie group countable at infinity. A continuous representation of \mathfrak{G} on a Hilbert space \mathcal{H} is a map $\pi: \mathfrak{G} \rightarrow \mathcal{L}(\mathcal{H})$, that satisfies:
 - ▶ π is a group homomorphism,
 - ▶ For every $\xi \in \mathcal{H}$, the function on \mathfrak{G} defined by $x \mapsto \pi(x)\xi$ is continuous.
- ▶ A representation π is called unitary if $\pi(x)$ is a unitary in $\mathcal{L}(\mathcal{H})$, for every $x \in \mathfrak{G}$.
- ▶ A representation is called irreducible if there are no non-trivial π -invariant closed subspaces of \mathcal{H} . In other words if $V \subseteq \mathcal{H}$ is a closed subspace, then $\pi(\mathfrak{G})V \subseteq V$ implies that either $V = 0$ or $V = \mathcal{H}$.

Lie Group Representations

- ▶ A vector $\xi \in \mathcal{H}$ is called smooth (for π) if the map $x \mapsto \pi(x)\xi$ is smooth on \mathfrak{G} .
- ▶ The space of all smooth vectors will be denoted by \mathcal{H}_∞ . This space is also called the Garding space of π .
- ▶ The Garding lemma tells us that \mathcal{H}_∞ is dense in \mathcal{H} and is π invariant.
- ▶ Topologize \mathcal{H}_∞ using the map $\xi \mapsto \pi(\cdot)\xi$. This map identifies the Garding space as a closed subspace of $C^\infty(\mathfrak{G}, \mathcal{H})$. This defines on \mathcal{H}_∞ a Frechet space topology, which is generally finer than the topology induced from \mathcal{H} .

Lie Group Representations

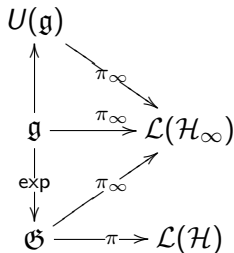
- ▶ Thus we get a representation of \mathfrak{G} on the Frechet space \mathcal{H}_∞ . We will denote this representation by π_∞ .
- ▶ Let \mathfrak{g} be the Lie algebra of \mathfrak{G} . Let $\mathfrak{g}_\mathbb{C}$ be the complexification of \mathfrak{g} and $U(\mathfrak{g})$ the universal enveloping algebra of $\mathfrak{g}_\mathbb{C}$.
- ▶ For every $\xi \in \mathcal{H}_\infty$ we have:

$$\forall \xi \in \mathcal{H}_\infty: \pi_\infty(X)\xi = \left. \frac{d}{dt} \right|_{t=0} \pi(\exp(tX))\xi.$$

- ▶ We can then extend π_∞ to a representation of $U(\mathfrak{g})$ on \mathcal{H}_∞ . We will denote this representation by π_∞ as well.

Lie Group Representations

A picture to describe what we have done so far:



Nuclear Spaces

- ▶ Let E, F be locally convex Hausdorff topological vector space.
- ▶ A continuous linear mapping $T: E \rightarrow F$ is called nuclear if there exist an equicontinuous sequence $\{x'_k\} \subset E'$, a sequence $\{y_k\} \subset B \subset F$, with B complete and a sequence $\{\lambda_k\} \subset \mathbb{C}$, such that $\sum_k |\lambda_k| < \infty$ and for every $x \in E$, we have:

$$T(x) = \sum_k \lambda_k \langle x'_k, x \rangle y_k.$$

- ▶ A space E is called nuclear if every continuous mapping $T: E \rightarrow B$ for any Banach space B is nuclear.

Nuclear Spaces

- ▶ If E is nuclear, then for every other locally convex Hausdorff topological vector space F the canonical map $E \otimes_{\pi} F \rightarrow E \otimes_{\epsilon} F$ is an isomorphism.
- ▶ This implies that essentially there is a unique way to define the tensor product of a locally convex Hausdorff TVS with a nuclear space.
- ▶ We are concerned with Frechet spaces and Banach spaces mostly. We have the following facts about them:
 - ▶ A Frechet space E is nuclear if and only if its dual is.
 - ▶ A normable space (in particular a Banach one) is nuclear if and only if it is finite dimensional.

Nuclear Spaces

- ▶ The abstract Schwartz Kernel Theorem states:

Theorem

Let E and F be a locally convex complete Hausdorff topological space. If in addition E is Frechet and nuclear then $\mathcal{L}_b(E, F)$ is complete and we have:

$$E' \widehat{\otimes} F \cong \mathcal{L}_b(E, F)$$

Frequency Domain Theory

- ▶ Recall that in the commutative case, if we plug $e^{\langle \lambda, t \rangle} u_0$, for $\lambda \in \mathbb{C}^2$, into the compatibility condition, we obtain an algebraic equation:

$$(\lambda_1 \sigma_2 - \lambda_2 \sigma_1 + i\gamma) u_0 = 0.$$

- ▶ One can also obtain this using the Fourier Transform of the function u .
- ▶ There is an analogue of the Fourier transform in the Lie group case.

Frequency Domain Theory

- ▶ We will now proceed to construct frequency domain theory for Lie algebra operator vessels.
- ▶ Let us assume that \mathfrak{g} is the Lie algebra of a second countable, countable at infinity Lie group \mathfrak{G} of type I .
- ▶ We will denote by $\widehat{\mathfrak{G}}$ the unitary dual of \mathfrak{G} . The collection of all equivalence classes of unitary irreducible representations.
- ▶ The idea is that $\pi \in \widehat{\mathfrak{G}}$ will be the frequency.

Frequency Domain Theory

- ▶ Applying the Fourier transform to a function f on \mathcal{G} we obtain an element of $\mathcal{L}(\mathcal{H}_\pi, \mathcal{H}_\pi \otimes \mathcal{E})$.
- ▶ For the following discussion to work we will restrict ourselves to the π -invariant subspace of smooth vectors, namely $\mathcal{H}_{\pi, \infty}$.
- ▶ We assume that $\mathcal{H}_{\pi, \infty}$ is nuclear and obtain from the Schwartz kernel theorem that:

$$\mathcal{L}(\mathcal{H}_{\pi, \infty}, \mathcal{H}_{\pi, \infty} \otimes \mathcal{E}) \cong \mathcal{H}'_{\pi, \infty} \otimes \mathcal{H}_{\pi, \infty} \otimes \mathcal{E}$$

Frequency Domain Theory

- ▶ The trajectories of the system are of the form:

$$u(g) = (\pi_\infty(g) \otimes I_{\mathcal{E}}) \mathbf{u}_0$$

$$x(g) = (\pi_\infty(g) \otimes I_{\mathcal{H}}) \mathbf{x}_0$$

$$y(g) = (\pi_\infty(g) \otimes I_{\mathcal{E}}) \mathbf{y}_0$$

- ▶ Here $\mathbf{u}_0, \mathbf{y}_0 \in \mathcal{H}'_{\pi, \infty} \otimes \mathcal{H}_{\pi, \infty} \otimes \mathcal{E}$ and $\mathbf{x}_0 \in \mathcal{H}'_{\pi, \infty} \otimes \mathcal{H}_{\pi, \infty} \otimes \mathcal{H}$.
- ▶ The frequency domain system is modified accordingly:

$$iX_k \mathbf{x} + [I_{\mathcal{H}_{\pi, \infty}} \otimes A_k] \mathbf{x} = [I_{\mathcal{H}_{\pi, \infty}} \otimes \Phi^* \sigma_k] \mathbf{u}$$

$$\mathbf{y} = \mathbf{u} - i[I_{\mathcal{H}_{\pi, \infty}} \otimes \Phi] \mathbf{x}$$

Frequency Domain Theory

- ▶ The i/o compatibility conditions then become:

$$(\pi_\infty(X_k) \otimes \sigma_j - \pi_\infty(X_j) \otimes \sigma_k + iI_{\mathcal{H}_{\pi,\infty}} \otimes \gamma_{jk})\mathbf{u} = 0$$

$$(\pi_\infty(X_k) \otimes \sigma_j - \pi_\infty(X_j) \otimes \sigma_k + iI_{\mathcal{H}_{\pi,\infty}} \otimes \gamma_{*jk})\mathbf{y} = 0$$

- ▶ Note that the operators $\pi_\infty(X_k)$ are continuous on $\mathcal{H}_{\pi,\infty}$.
- ▶ Using the Fourier transform, one can translate between the frequency domain and time domain systems of a Lie algebra operator vessel.

Frequency Domain Theory

- ▶ We then define the joint transfer function of the vessel by:

$$S(\pi) = I_{\mathcal{H}_{\pi, \infty}} \otimes I_{\mathcal{E}} - i (I_{\mathcal{H}_{\pi, \infty}} \otimes \Phi) R(X, \pi) (I_{\mathcal{H}_{\pi, \infty}} \otimes \Phi^* \sigma(X)).$$

Where $R(X, \pi) = (i\pi_{\infty}(X) \otimes I_{\mathcal{H}} + I_{\mathcal{H}_{\pi, \infty}} \otimes \rho(X))^{-1}$.

- ▶ We denote the space of solution of the input compatibility conditions by $\mathcal{E}(\pi) \subset \mathcal{L}(\mathcal{H}_{\pi, \infty}, \mathcal{H}_{\pi, \infty} \otimes \mathcal{E})$. Similarly we denote by $\mathcal{E}_*(\pi)$ the space of solutions of the output compatibility conditions.
- ▶ $S(\pi)$ maps the space $\mathcal{E}(\pi)$ into $\mathcal{E}_*(\pi)$.

Frequency Domain Theory

- ▶ The next example is the $ax + b$ group. This is the smallest non-commutative group. It is the group of affine transformations of the plane.
- ▶ The Plancherel measure on the dual is the counting measure on two points representing the only two classes of infinite-dimensional irreducible representations, π_+ , a representation on $L^2(\mathbb{R}_{>0}, ds/s)$, and π_- , a representation on $L^2(\mathbb{R}_{<0}, ds/s)$.
- ▶ Then using the Schwartz Kernel Theorem again we obtain that the input compatibility conditions are differential equations for vector-valued distributions on \mathbb{R}^2 .

Frequency Domain

- ▶ After a complexification process we obtain the following PDEs:

$$\frac{1}{2}\sigma_2 K(z, w) + z\sigma_2 \frac{\partial}{\partial z} K(z, w) - 2\pi iz\sigma_1 K(z, w) + i\gamma K(z, w) = 0$$

$$\frac{1}{2}\sigma_2 K(z, w) + z\sigma_2 \frac{\partial}{\partial z} K(z, w) - 2\pi iz\sigma_1 K(z, w) + i\gamma_* K(z, w) = 0$$

- ▶ The joint transfer function is then:

$$S(\pi) = I_{\mathcal{E}} - i\Phi(A_2 - 2\pi tI_{\mathcal{H}})^{-1}\Phi^*\sigma_2$$

This is an isomonodromic mapping that maps solutions of the above PDE's one to another:

Taylor Joint Spectrum

- ▶ We will describe the spectrum of a Lie algebra representation following [Taylor, 1972].
- ▶ Let \mathfrak{g} be a Lie algebra of dimension n , $\mathfrak{g}_{\mathbb{C}}$ its complexification and $U(\mathfrak{g})$ the universal enveloping algebra of $\mathfrak{g}_{\mathbb{C}}$.
- ▶ Every representation of \mathfrak{g} on a locally convex space X can be extended to a representation of $\mathfrak{g}_{\mathbb{C}}$ and $U(\mathfrak{g})$ on X . Thus one gets a structure of a left $U(\mathfrak{g})$ module on X .
- ▶ On the other hand every left $U(\mathfrak{g})$ module defines a representation of \mathfrak{g} .

Taylor Joint Spectrum

- ▶ Consider the augmented Koszul complex, that is a free resolution of $U(\mathfrak{g})$ as a topological bimodule over itself:

$$0 \longleftarrow U(\mathfrak{g}) \longleftarrow U(\mathfrak{g}) \widehat{\otimes} U(\mathfrak{g}) \longleftarrow$$

$$\longleftarrow U(\mathfrak{g}) \widehat{\otimes} U(\mathfrak{g}) \otimes \mathfrak{g} \longleftarrow U(\mathfrak{g}) \widehat{\otimes} U(\mathfrak{g}) \otimes \wedge^2 \mathfrak{g} \longleftarrow$$





$$\longleftarrow \dots \longleftarrow U(\mathfrak{g}) \widehat{\otimes} U(\mathfrak{g}) \otimes \wedge^n \mathfrak{g}$$

Taylor Joint Spectrum

- ▶ Let M be a $U(\mathfrak{g})$ topological bimodule. One can then compute the Hochschild homology of M , $H_k(U(\mathfrak{g}), M)$, by tensoring M (in the category of $U(\mathfrak{g})$ topological bimodules) with the above Koszul complex.
- ▶ Now we define the resolvent set of \mathfrak{g} to be the topological $U(\mathfrak{g})$ -bimodules, such that the $H_k(U(\mathfrak{g}), M)$ vanish for all $k > 0$.
- ▶ How is this related to Lie algebra operator vessels?

Taylor Joint Spectrum

- ▶ Let \mathfrak{G} be a Lie group as in the previous section. Let $\pi \in \widehat{\mathfrak{G}}$, be an irreducible unitary representation of \mathfrak{G} on \mathcal{H}_π . Finally let $\mathcal{H}_{\pi,\infty}$ be the space of smooth vectors of π .
- ▶ Let ρ be a representation of \mathfrak{g} on a Hilbert space \mathcal{H} . it then extends uniquely to a representation of $\mathfrak{g}_\mathbb{C}$ on \mathcal{H} , that we will also call ρ . Consider $M_\pi = \mathcal{H}_{\pi,\infty} \otimes \mathcal{H}$ as a $U(\mathfrak{g})$ bimodule with the obvious action.
- ▶ **Question:** Assume that the imaginary part of $\rho(X)$ is compact, for every $X \in \mathfrak{g}$. Is it true then that the resolvent in the definition of the joint transfer function of the vessel defined by ρ exists if and only if M_π is in the resolvent set of $U(\mathfrak{g})$?

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Thank You!