



## On J-Symmetric Restricted Shifts

Paul A. Fuhrmann

*Proceedings of the American Mathematical Society*, Vol. 51, No. 2. (Sep., 1975), pp. 421-426.

Stable URL:

<http://links.jstor.org/sici?sici=0002-9939%28197509%2951%3A2%3C421%3AORS%3E2.0.CO%3B2-T>

*Proceedings of the American Mathematical Society* is currently published by American Mathematical Society.

---

Your use of the JSTOR archive indicates your acceptance of JSTOR's Terms and Conditions of Use, available at <http://www.jstor.org/about/terms.html>. JSTOR's Terms and Conditions of Use provides, in part, that unless you have obtained prior permission, you may not download an entire issue of a journal or multiple copies of articles, and you may use content in the JSTOR archive only for your personal, non-commercial use.

Please contact the publisher regarding any further use of this work. Publisher contact information may be obtained at <http://www.jstor.org/journals/ams.html>.

Each copy of any part of a JSTOR transmission must contain the same copyright notice that appears on the screen or printed page of such transmission.

---

JSTOR is an independent not-for-profit organization dedicated to and preserving a digital archive of scholarly journals. For more information regarding JSTOR, please contact [support@jstor.org](mailto:support@jstor.org).

## ON $J$ -SYMMETRIC RESTRICTED SHIFTS<sup>1</sup>

PAUL A. FUHRMANN

ABSTRACT. The restricted shift operators in proper left invariant subspaces of  $H^2$  that are  $J$ -symmetric are characterized and the signature of the corresponding operator  $J$  is determined.

The fact that left shifts in vectorial  $H^2$  spaces restricted to their left invariant subspaces can serve as models for the most general contractions in Hilbert space is well known [7], [13]. Thus the structural analysis of such operators contributes to the general structure theory. Of course in the general case we will have to consider shifts of infinite multiplicity.

For shifts of finite multiplicity more is known [7] and they have, besides their intrinsic mathematical interest, considerable importance due to their use as models for the generators in internal descriptions of linear time invariant dynamical systems. In fact most of finite dimensional system theory can be done, in an elegant way, using operator theoretic methods with an emphasis on Hankel and shift operators. We quote a few papers [1], [5], [6], [8] that can serve as a guide to the interested reader. This paper itself has been motivated by problems of system theory where the signature of the operator  $J$  is related to the residues of the transfer function of the system.

In a Hilbert space  $H$  we consider a bounded operator  $J$  satisfying  $J = J^* = J^{-1}$ . This implies there exist two orthogonal projections  $P_+$  and  $P_-$  for which  $I = P_+ + P_-$ ,  $J = P_+ - P_-$  and  $P_+P_- = 0$ . Thus if  $H_{\pm} = P_{\pm}H$  then clearly  $H_{\pm} = \{x \in H \mid Jx = \pm x\}$ . A bounded operator  $A$  is called  $J$ -symmetric if  $A = JA^*J$ . The  $J$ -symmetric operators have been widely studied and [10], [11] are some references to the literature. Our interest will be in a very special class of  $J$ -symmetric operators.

---

Presented to the Society, February 18, 1974; received by the editors March 5, 1974 and, in revised form, May 23, 1974.

AMS (MOS) subject classifications (1970). Primary 47A45; Secondary 47B50.

Key words and phrases. Invariant subspaces, restricted shifts,  $J$ -symmetric operators.

<sup>1</sup>This work was supported by the US Office of Naval Research under the Joint Services Electronics Program by Contract N00014-67-A-0298-0006.

Let  $S$  be the right shift in  $H^2$  [9] i.e. the operator defined by  $(Sf)(z) = zf(z)$ . Let  $K$  be a proper left invariant subspace of  $H^2$ , that is, a subspace of  $H^2$  invariant under  $S^*$ . By a theorem of Beurling [9],  $K^\perp = \phi H^2$  for some inner function  $\phi$ . Let  $P_{\{\phi H^2\}^\perp}$  be the orthogonal projection of  $H^2$  onto  $K = \{\phi H^2\}^\perp$ . We define an operator  $T_\phi$  in  $\{\phi H^2\}^\perp$  by  $T_\phi f = P_{\{\phi H^2\}^\perp} S f$  for all  $f \in \{\phi H^2\}^\perp$  and then we have  $T_\phi^* = S^*|_{\{\phi H^2\}^\perp}$ . Since  $T_\phi$  is completely determined by the inner function  $\phi$  we want to get the relation between the analytic properties of  $\phi$  and the  $J$ -symmetry of  $T_\phi$ . Inner functions are determined only up to a constant factor of modulus one, thus we will normalize the inner functions by requiring their first nonvanishing Taylor coefficient to be positive. For every  $a \in H^\infty$  we let  $\tilde{a}(z) = \overline{a(\bar{z})}$ . Thus  $a = \tilde{a}$  if and only if all coefficients in a power series expansion of  $a$  at zero are real.

For every inner function  $\phi$  we define  $\tau_\phi$  on  $\{\phi H^2\}^\perp$  by

$$(1) \quad (\tau_\phi f)(e^{it}) = e^{-it} \tilde{\phi}(e^{it}) f(e^{-it}),$$

$\tau_\phi$  is a unitary map of  $\{\phi H^2\}^\perp$  onto  $\{\tilde{\phi} H^2\}^\perp$  for which  $T_{\tilde{\phi}} \tau_\phi = \tau_\phi T_\phi^*$  [4]. Clearly  $\tau_\phi^{-1} = \tau_\phi^* = \tau_{\tilde{\phi}}$ . Thus if  $\phi = \tilde{\phi}$  then  $\tau_\phi$  is a  $J$  operator in  $\{\phi H^2\}^\perp$ .

**Theorem 1.**  *$T$  is a  $J$ -symmetric operator if and only if  $\phi = \tilde{\phi}$ . The corresponding  $J$  is given by  $J = \pm \tau_\phi$ .*

Before proceeding with the proof of the theorem we establish

**Lemma 1.** *The only unitary operators in  $\{\phi H^2\}^\perp$  commuting with  $T_\phi$  are multiplications by constants of modulus one.*

**Proof.** One part is trivial. So let us assume  $V$  is unitary in  $\{\phi H^2\}^\perp$  and  $T_\phi V = VT_\phi$ . By Sarason's commutant theorem [12] we have  $V = v(T)$  for some  $v \in H^\infty$  satisfying  $\|v\|_\infty = 1$ . Since  $V$  is unitary we have  $|v(e^{it})| = 1$ , i.e. either  $v$  is inner or a constant of modulus one. Assume  $v$  is inner, then  $v^n f \perp \phi H^2$  for all  $f \in \{\phi H^2\}^\perp$  and all  $n \geq 0$ . Now  $f \in \{\phi H^2\}^\perp$  if and only if on the unit circle  $f$  has the factorization  $f(e^{it}) = e^{-it} \phi(e^{it}) \overline{h(e^{it})}$  for some  $h \in H^2$  [3]. Thus our assumption implies  $\bar{z} v^n \bar{h} \perp H^2$  for all  $n \geq 0$ . This implies  $h \in \bigcap_{n=0}^\infty v^n H^2 = \{0\}$ . Thus  $h = 0$  and hence  $f = 0$ . Since  $f$  is arbitrary and  $\phi$  nontrivial, this is impossible. Thus  $v$  is a constant of modulus one.

**Proof of Theorem 1.** If  $\phi = \tilde{\phi}$  it follows from the remark following (1) that  $T_\phi$  is  $J$ -symmetric with respect to  $\tau_\phi$ . Conversely let us assume  $T_\phi$

to be  $J$ -symmetric with respect to an arbitrary  $J$ . Hence  $T_\phi^* = JT_\phi J$ . This implies that the minimal inner functions [13] of  $T_\phi$  and  $T_\phi^*$  are the same, therefore  $\phi = \tilde{\phi}$ . This in turn implies, as above, that  $T_\phi = \tau_\phi T_\phi^* \tau_\phi$  and hence  $T_\phi = \tau_\phi (JT_\phi J) \tau_\phi = (\tau_\phi J) T_\phi (J\tau_\phi)$  or  $T_\phi (J\tau_\phi) = (J\tau_\phi) T_\phi$ . Thus  $J\tau_\phi$ , which is unitary, commutes with  $T_\phi$  and hence, by Lemma 1,  $J\tau_\phi = \alpha$  with  $\alpha$  a constant of modulus one. Multiplying on the left by  $J$  we have  $\tau_\phi = \alpha J$  and as  $\tau_\phi$  and  $J$  are selfadjoint it follows that  $\alpha$  is real. Hence  $\alpha = \pm 1$ .

Next we want to find the dimensions of  $K_+ = \{f \in K \mid \tau_\phi f = f\}$  and  $K_- = \{f \in K \mid \tau_\phi f = -f\}$ . Let us denote by  $k_0$  and  $K_0$  the functions defined by  $k_0 = P_{\{\phi H^2\}^\perp} 1$  and  $K_0 = P_{\{\phi H^2\}^\perp} \bar{z}\phi$  [2]. Since

$$T_\phi^n k_0 = P_{\{\phi H^2\}^\perp} z^n k_0 = P_{\{\phi H^2\}^\perp} z^n,$$

we have the set of vectors  $\{T_\phi^n k_0 \mid n \geq 0\}$  spanning  $\{\phi H^2\}^\perp$  or equivalently  $k_0$  is a cyclic vector for  $T_\phi$ . Now  $\tau_\phi k_0 = \tilde{K}_0$  and  $\tau_\phi K_0 = \tilde{k}_0$  which implies that  $\tau_\phi T_\phi^* K_0 = T_\phi^* \tilde{k}_0$  and hence, by the fact that  $\tau_\phi$  is unitary,  $K_0$  is a cyclic vector for  $T_\phi^*$ , or  $\{T_\phi^{*n} K_0 \mid n \geq 0\}$  also span  $\{\phi H^2\}^\perp$ .

**Theorem 2.** *Let  $\phi$  be inner satisfying  $\phi = \tilde{\phi}$ .*

a.  $K_+ = \text{Clos}\{\alpha(T)k_0 + \alpha(T^*)K_0 \mid \alpha \in H^\infty\},$

$K_- = \text{Clos}\{\alpha(T)k_0 - \alpha(T^*)K_0 \mid \alpha \in H^\infty\}.$

b.  $K_+$  is finite dimensional if and only if  $\phi$  is a finite Blaschke product. The same holds for  $K_-$ .

c. If  $\phi$  is a finite Blaschke product of  $n$  factors then  $\dim K_+ = [(n+1)/2]$  and  $\dim K_- = [n/2]$ .

**Proof.** a. Since  $\phi = \tilde{\phi}$  we have  $\tau_\phi T_\phi^* = T_\phi \tau_\phi$  and in this case  $\tau_\phi k_0 = K_0$  and  $\tau_\phi K_0 = k_0$ . Thus

$$\tau_\phi \{\alpha(T)k_0 + (T^*)K_0\} = \{\alpha(T^*)K_0 + \alpha(T)k_0\}$$

and this implies that for all  $f$  in  $K_+ = \text{Clos}\{\alpha(T)k_0 + \alpha(T^*)K_0 \mid \alpha \in H^\infty\}$  we have  $\tau_\phi f = f$ . Similarly  $\tau_\phi g = -g$  for all  $g$  in  $K_- = \text{Clos}\{\alpha(T)k_0 - \alpha(T^*)K_0 \mid \alpha \in H^\infty\}$ . Clearly those subspaces are orthogonal for if  $f \in K_+$  and  $g \in K_-$  then

$$(f, g) = (\tau_\phi f, g) = (f, \tau_\phi g) = (f, -g) = -(f, g)$$

and hence  $(f, g) = 0$ . The direct sum of these two subspaces contains all vectors of the form  $T_\phi^n k_0, n \geq 0$ , and hence is all of  $\{\phi H^2\}^\perp$ .

b.  $\{\phi H^2\}^\perp$  is finite dimensional if and only if  $\phi$  is a Blaschke product with a finite number of factors. Thus in that case  $K_+$  and  $K_-$  are finite dimensional. Assume now  $K_-$  is finite dimensional, in this case only a finite number of functions of the form  $T^k k_0 - T^{*k} K_0$  are linearly independent. Thus for some integer  $n$  and, not all zero,  $\alpha_j$  we have  $\sum_{j=1}^n \alpha_j (T^j k_0 - T^{*j} K_0) = 0$ . The functions  $T^k k_0$  and  $T^{*j} K_0$  are easily expressible in terms of the inner function  $\phi$  and in general we have

$$(T^j k_0)(z) = z^j - \phi(z) \sum_{k=0}^j \bar{\phi}_k z^{j-k}$$

and

$$(T^{*j} K_0)(z) = z^{-(j+1)} \left[ \phi(z) - \sum_{k=0}^j \phi_k z^k \right].$$

Thus if  $\sum_{j=0}^n \alpha_j (T^j k_0 - T^{*j} K_0) = 0$  it follows, dropping bars as  $\phi = \bar{\phi}$ , that

$$\begin{aligned} \sum_{j=0}^n \alpha_j z^j - \phi(z) \sum_{j=0}^n \alpha_j \sum_{k=0}^j \phi_k z^{j-k} \\ = z^{-(n+1)} \phi(z) \sum_{j=0}^n \alpha_j z^{h-j} - z^{-(n+1)} \sum_{j=0}^n \alpha_j z^{n-j} \sum_{k=0}^j \phi_k z^k. \end{aligned}$$

Now

$$\sum_{j=0}^n \alpha_j \sum_{k=0}^j \phi_k z^{j-k} = \sum_{j=0}^n \left[ \sum_{k=j}^n \alpha_k \phi_{k-j} \right] z^j$$

and

$$\sum_{j=0}^n \alpha_j z^{n-j} \sum_{k=0}^j \phi_k z^k = \sum_{j=0}^n \left[ \sum_{k=0}^j \alpha_k \phi_{n+k-j} \right] z^j.$$

Therefore  $p(z) = \phi(z)q(z)$  with  $p$  and  $q$  being  $(2n+1)$ -degree polynomials given by

$$p(z) = \sum_{j=0}^n \alpha_j z^{n+j+1} + \sum_{j=0}^n \left[ \sum_{k=j}^n \alpha_k \phi_{k-j} \right] z^j$$

and

$$q(z) = \sum_{j=0}^n \alpha_j z^{n-j} + \sum_{j=0}^n \left[ \sum_{k=j}^n \alpha_k \phi_{k-j} \right] z^{j+n+1}.$$

It is clear from the above that if  $p_i$  and  $q_i$  are the coefficients of  $p$  and  $q$  respectively then  $p_i = q_{2n+1-i}$ . It follows that  $\phi$  is a Blaschke product of at most  $2n + 1$  factors. The same result follows from the assumption  $\dim K_+ = n$ .

c. Let us assume now that  $\dim K = n$ , i.e.  $\phi$  is a Blaschke product of  $n$  factors. Let the zeroes of  $\phi$ , repeated according to multiplicities, be  $\lambda_i, i = 1, \dots, n$ . Thus  $\phi(z) = \prod_{i=1}^n (z - \lambda_i)/(1 - \bar{\lambda}_i z)$ . Let  $m(z) = \prod_{i=1}^n (z - \lambda_i)$ . Thus  $m$  is the minimal polynomial of  $T_\phi$ , and as  $m = \tilde{m}$ , also of  $T_\phi^*$ . Clearly we have  $\phi(e^{it}) = m(e^{it})/e^{int}m(e^{-it})$ . Let  $\tau = \tau_{z^n}$  be defined in  $\{z^n H^2\}^\perp$  by (1). We will show that  $\tau_\phi$  and  $\tau$  are similar and hence their signatures are the same. To this end we define a map  $\Phi$  on  $\{\phi H^2\}^\perp$

$$(2) \quad (\Phi f)(e^{it}) = e^{int}m(e^{-it})f(e^{it}).$$

Let  $p(e^{it}) = e^{int}m(e^{-it})$ ; then  $p$  is a polynomial of degree  $n$  and hence  $\Phi f = pf$  and  $pf \in H^2$  for all  $f \in \{\phi H^2\}^\perp$ . Now  $m$  being the minimal polynomial of  $T_\phi^*$  implies  $\bar{m}f \perp H^2$  for all  $f \in \{\phi H^2\}^\perp$ . Thus  $z^n \bar{m}f = pf$  is a polynomial of degree  $\leq n - 1$ . So  $\Phi$  maps  $\{\phi H^2\}^\perp$  into  $\{z^n H^2\}^\perp$ . Since  $p \neq 0$   $\Phi$  is 1-1 and hence an invertible map of  $\{\phi H^2\}^\perp$  onto  $\{z^n H^2\}^\perp$ . Now

$$\begin{aligned} (\Phi \tau_\phi f)(e^{it}) &= p(e^{it})(\tau_\phi f)(e^{it}) = p(e^{it})\phi(e^{it})f(e^{-it}) \\ &= e^{-it}e^{int}m(e^{-it})\phi(e^{it})f(e^{-it}) = e^{-it}m(e^{it})f(e^{-it}) \\ &= e^{-it} \cdot e^{int} \cdot e^{-int}m(e^{it})f(e^{-it}) = (\tau \Phi f)(e^{it}). \end{aligned}$$

Hence  $\Phi \tau_\phi = \tau \Phi$  and this means that  $\tau_\phi$  and  $\tau$  are similar. In particular  $\tau_\phi f = f$  if and only if  $\tau \Phi f = \Phi f$  and  $\tau_\phi f = -f$  if and only if  $\tau(\Phi f) = -(\Phi f)$ . Thus the signatures of  $\tau$  and  $\tau_\phi$  are the same. In terms of the natural orthonormal basis of  $\{z^n H^2\}^\perp$  consisting of the functions  $1, z, \dots, z^{n-1}$ ,  $\tau$  has a matrix representation given by  $T = (t_{ij}), t_{ij} = 0$  if  $i + j \neq n$  and 1 if  $i + j = n$ . Clearly  $\text{tr}(T) = 0$  if  $n$  is even and 1 if it is odd, hence the result.

**Corollary.** *Let  $\phi$  be inner.  $T_\phi$  defined by (1) in  $\{\phi H^2\}^\perp$  is selfadjoint if and only if  $\phi(z) = \alpha(z - \lambda)(1 - \lambda z)^{-1}$  for some real  $\lambda, |\lambda| < 1$ , and  $\alpha$  of modulus one.*

REFERENCES

1. J. S. Baras and R. W. Brockett,  $H^2$ -functions and infinite dimensional realization theory, SIAM J. Control 13 (1975), 221-241.
2. D. N. Clark, One dimensional perturbations of restricted shifts, J. Analyse Math. 25 (1972), 169-191. MR 46 #692.

3. R. G. Douglas, H. S. Shapiro and A. L. Shields, *Cyclic vectors and invariant subspaces for the backward shift operator*, Ann. Inst. Fourier (Grenoble) 20 (1970), fosc. 1, 37–76. MR 42 # 5088.
4. P. A. Fuhrmann, *On the Corona theorem and its applications to spectral problems in Hilbert space*, Trans. Amer. Math. Soc. 132 (1968), 55–66. MR 36 # 5751.
5. ———, *On realization of linear systems and applications to some questions of stability*, Math. Systems Theory 8 (1975), 132–141.
6. ———, *Realization theory in Hilbert space for a class of transfer functions*, J. Functional Analysis 18 (1975), 338–349.
7. H. Helson, *Lectures on invariant subspaces*, Academic Press, New York and London, 1964. MR 30 # 1409.
8. J. W. Helton, *Discrete time systems operator models and scattering theory*, J. Functional Analysis 16 (1974), 15–38.
9. K. Hoffman, *Banach spaces of analytic functions*, Prentice-Hall Ser. in Modern Analysis, Prentice-Hall, Englewood Cliffs, N. J., 1962. MR 24 # A2844.
10. I. S. Iohvidov and M. G. Kreĭn, *Spectral theory of operators in spaces with indefinite metric*. I, Trudy Moscov. Mat. Obšč. 5 (1956), 367–432; English transl., Amer. Math. Soc. Transl. (2) 13 (1960), 105–175. MR 18, 320; 22 # 3983.
11. M. G. Kreĭn, *Introduction to the geometry of indefinite  $J$ -spaces and to the theory of operators in those spaces*, Second Math. Summer School, part 1, Naukova Dumka, Kiev, 1965, pp. 15–92; English transl., Amer. Math. Soc. Transl. (2) 93 (1970), 103–176. MR 33 # 574; 42 # 4.
12. D. Sarason, *Generalized interpolation in  $H^\infty$* , Trans. Amer. Math. Soc. 127 (1967), 179–203. MR 34 # 8193.
13. B. Sz.-Nagy and C. Foiaş, *Analyse harmonique des opérateurs de l'espace de Hilbert*, Masson, Paris; Akad. Kiadó, Budapest, 1967; English rev. transl., North-Holland, Amsterdam; American Elsevier, New York; Akad. Kiadó, Budapest, 1970. MR 37 # 778; 43 # 947.

DIVISION OF ENGINEERING AND APPLIED PHYSICS, HARVARD UNIVERSITY,  
CAMBRIDGE, MASSACHUSETTS 02138

*Current address:* Department of Mathematics, Ben Gurion University of the  
Negev, Beer Sheva, Israel